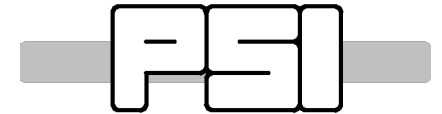
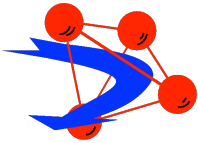


# Determination of long range antiferromagnetic order by powder neutron diffraction

Vladimir Pomjakushin

*Laboratory for Neutron Scattering and Imaging, LNS, Paul Scherrer Institute*



During the practicum we will try to reproduce the one of the neutron diffraction experiments performed during 1946-1951 for which C.G. Shull was honoured with the Nobel Prize in 1994. We will perform neutron diffraction experiment with MnS using powder diffractometer HRPT/SINQ. From the analysis of the nuclear and magnetic Bragg peak intensities and positions we will verify the crystal and magnetic structures of manganese sulfide MnS and determine the size of the magnetic moment on manganese.

# Determination of long range antiferromagnetic order by powder neutron diffraction

- Historical introduction (4)
- Some experimental details of neutron diffraction experiment at HRPT/SINQ (10)
- Reminder on nuclear and magnetic neutron structure factors (6)
- Practicum problems (6)

# 1994 Nobel Prize in Physics

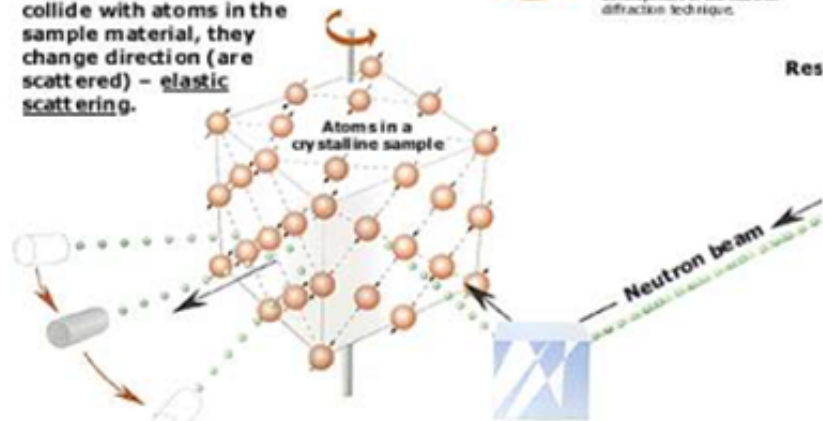
Clifford G. Shull  
1915 – 2001, USA



Clifford G. Shull, MIT, Cambridge, Massachusetts, USA, receives one half of the 1994 Nobel Prize in Physics for development of the neutron diffraction technique.

Neutrons show where atoms are

When the neutrons collide with atoms in the sample material, they change direction (are scattered) – elastic scattering.



Detectors record the directions of the neutrons and a diffraction pattern is obtained. The pattern shows the positions of the atoms relative to one another.

Crystal that sorts and forwards neutrons of a certain wavelength (energy) – monochromatized neutrons

Res

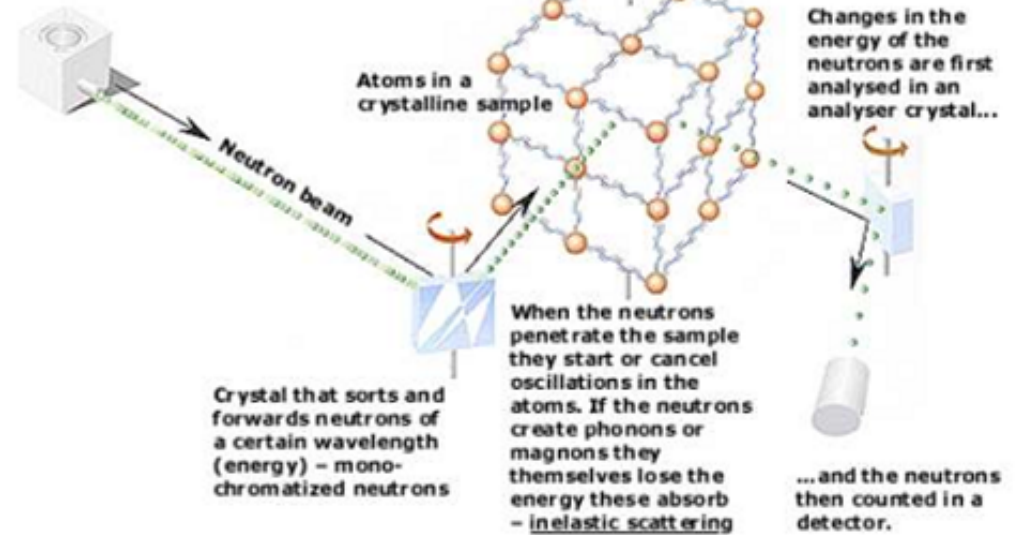
Bertram N. Brockhouse  
1918 – 2003, Canada



Bertram N. Brockhouse, McMaster University, Hamilton, Ontario, Canada, receives one half of the 1994 Nobel Prize in Physics for the development of neutron spectroscopy.

Neutrons show what atoms do

3-axis spectrometer with rotatable crystals and rotatable sample



Crystal that sorts and forwards neutrons of a certain wavelength (energy) – monochromatized neutrons

When the neutrons penetrate the sample they start or cancel oscillations in the atoms. If the neutrons create phonons or magnons they themselves lose the energy these absorb – inelastic scattering

Changes in the energy of the neutrons are first analysed in an analyser crystal...

...and the neutrons then counted in a detector.

# Neutron Diffraction by Paramagnetic and Antiferromagnetic Substances

C. G. SHULL, W. A. STRAUER, AND E. O. WOLLAN  
*Oak Ridge National Laboratory, Oak Ridge, Tennessee*

(Received March 2, 1951)

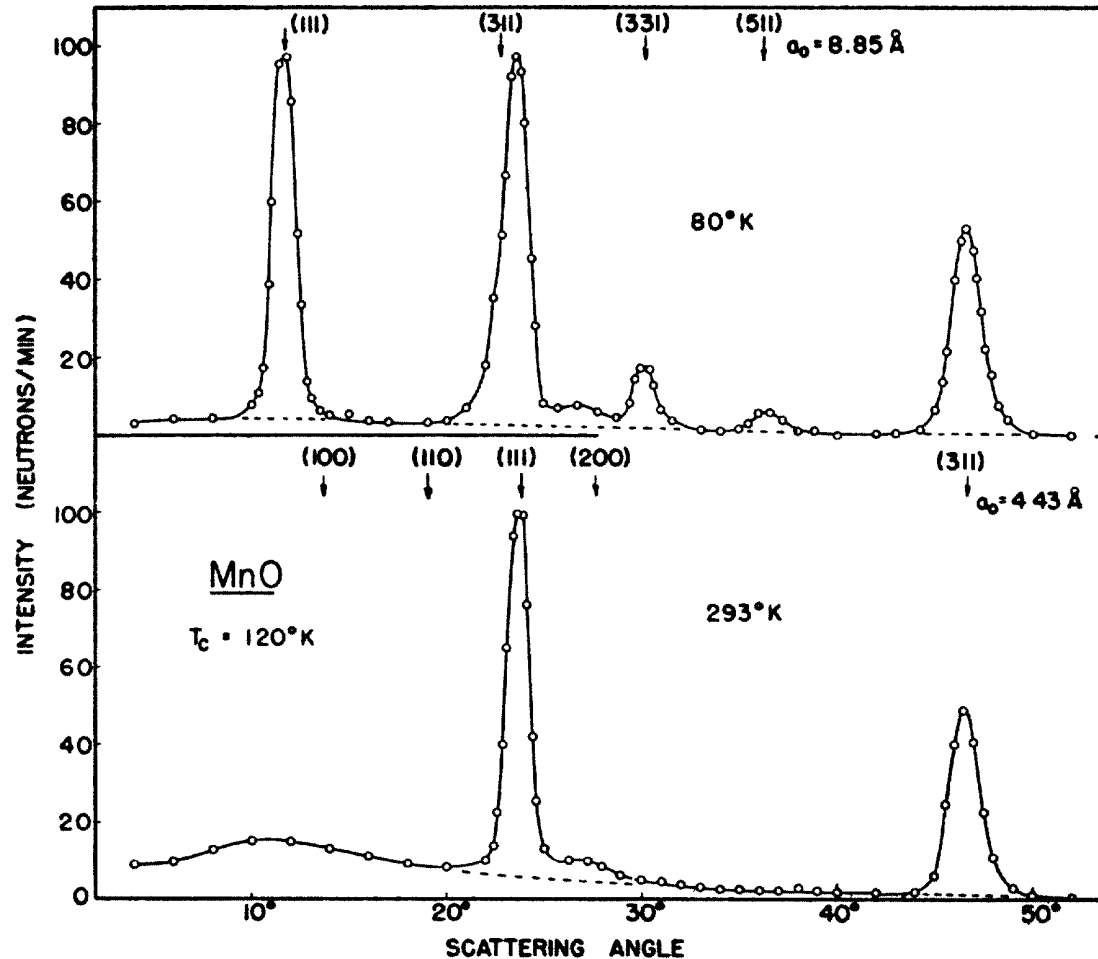


FIG. 4. Neutron diffraction patterns for MnO taken at liquid nitrogen and room temperatures. The patterns have been corrected for the various forms of extraneous, diffuse scattering mentioned in the text. Four extra antiferromagnetic reflections are to be noticed in the low temperature pattern.

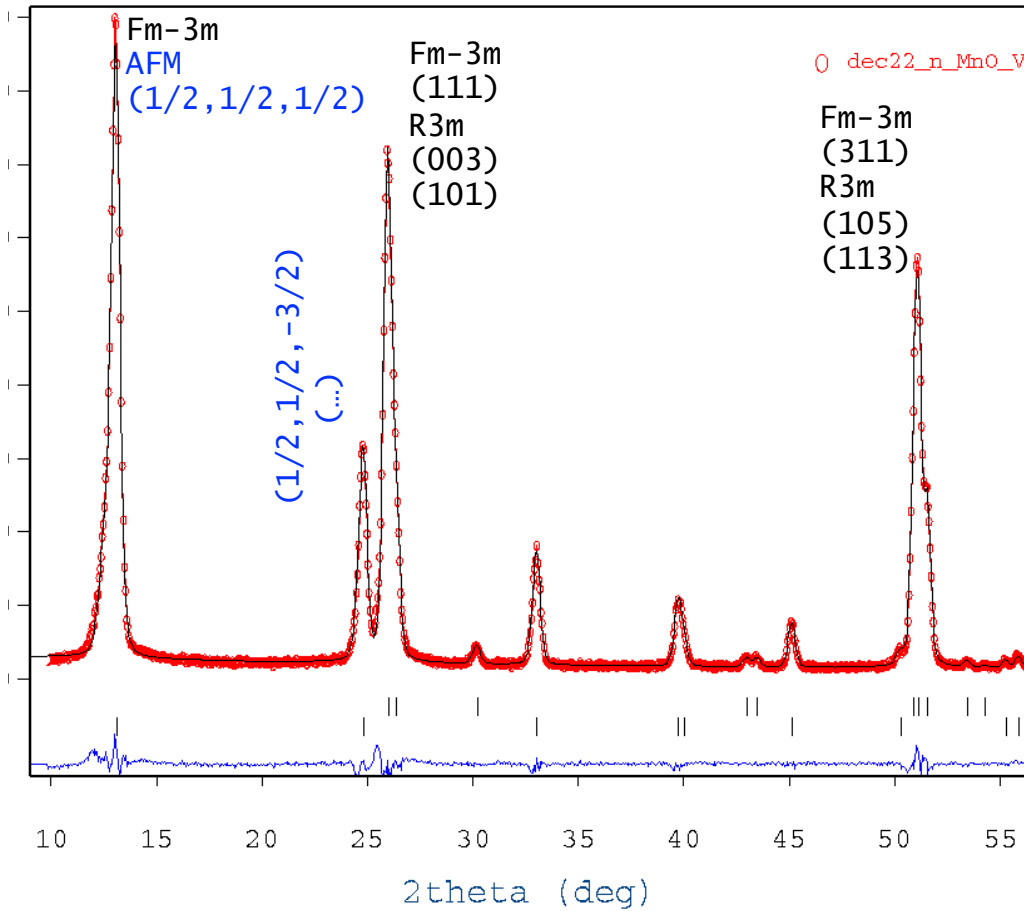


# HRPT/SINQ nowadays

$\lambda=1.15\text{\AA}$ , MnO @ 2K.

Rhombohedral distortions are explicitly seen

R-3m and k=003/2



# reaction by Paramagnetic and Antiferromagnetic

$\lambda=1.057\text{\AA}$

C. G. SHULL, W. A. STRAUER, AND E. O. WOLLAN  
 Oak Ridge National Laboratory, Oak Ridge, Tennessee

(Received March 2, 1951)

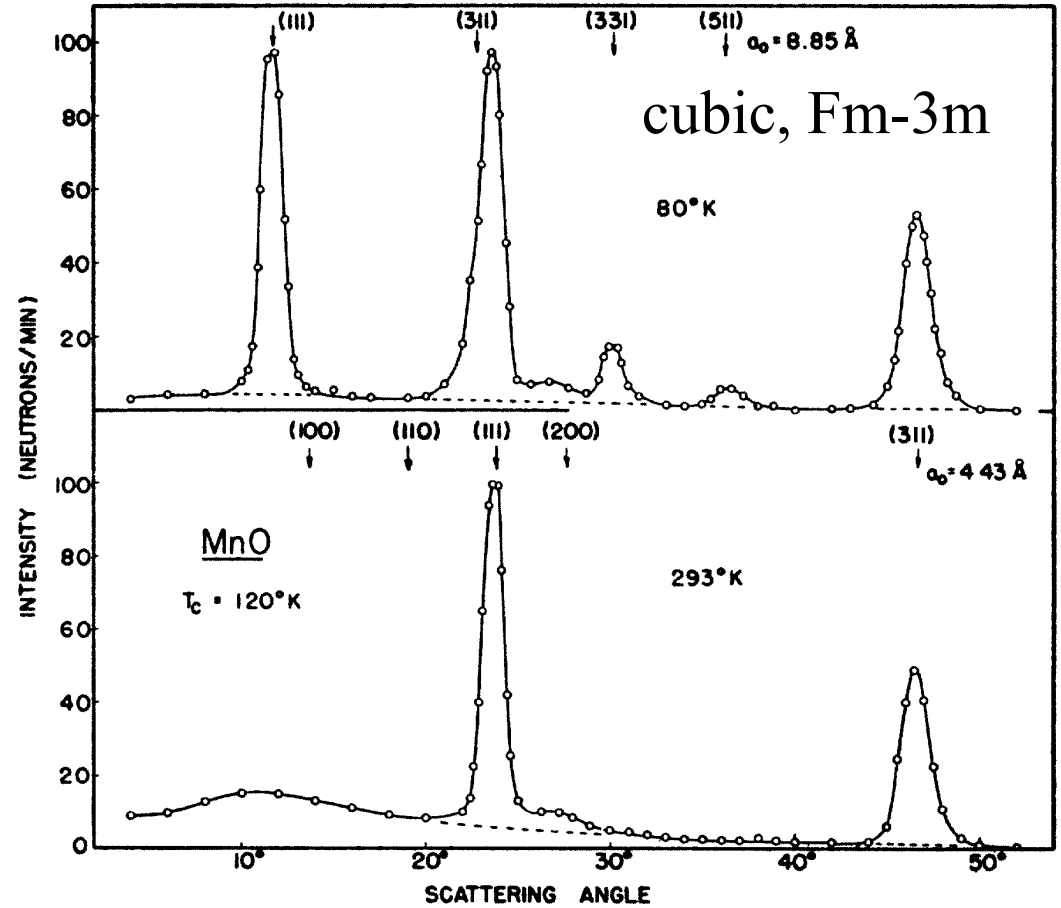
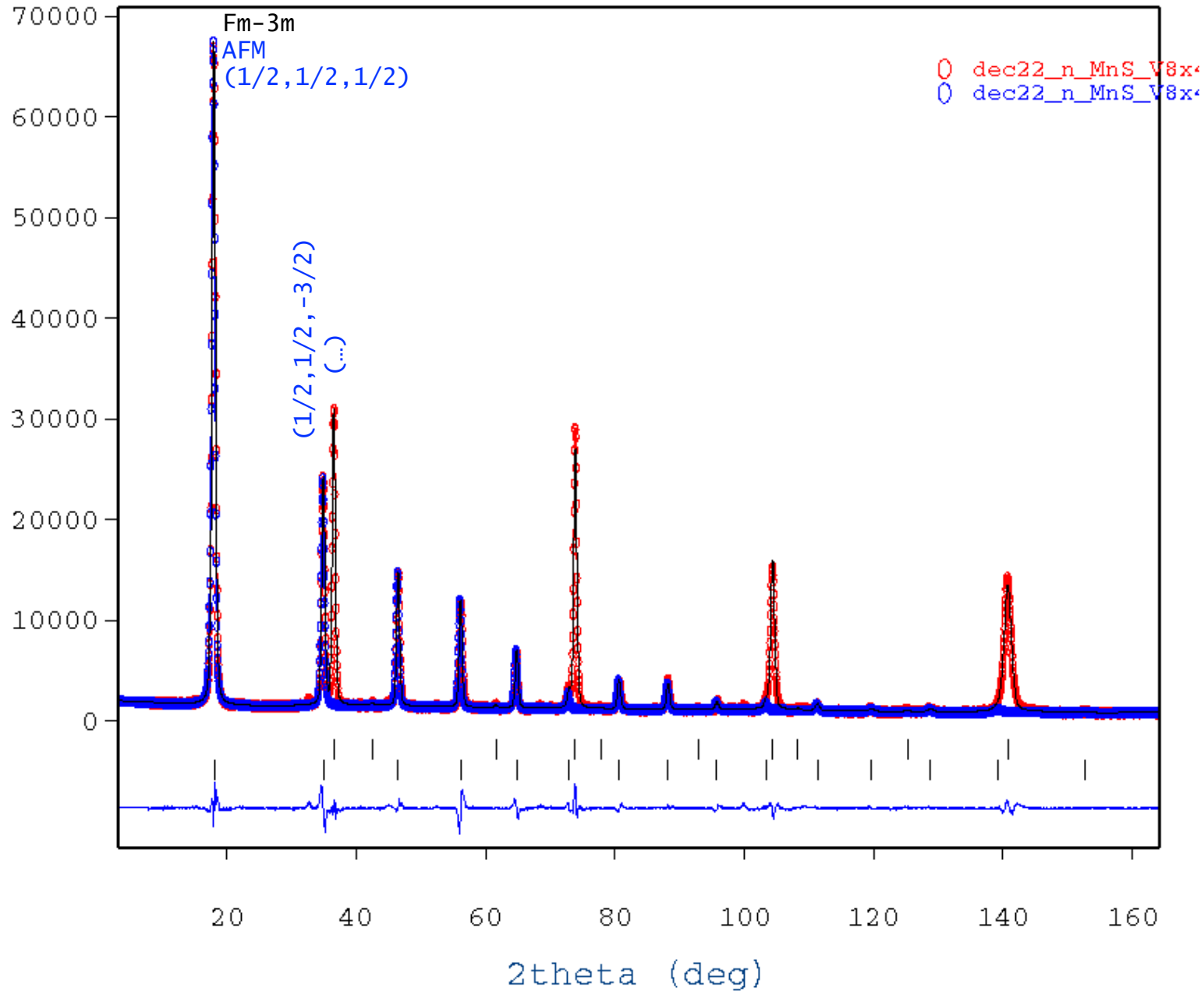


FIG. 4. Neutron diffraction patterns for MnO taken at liquid nitrogen and room temperatures. The patterns have been corrected for the various forms of extraneous, diffuse scattering mentioned in the text. Four extra antiferromagnetic reflections are to be noticed in the low temperature pattern.

# HRPT, 1.9Å, MnS @ T=2K, pseudo-cubic Fm-3m

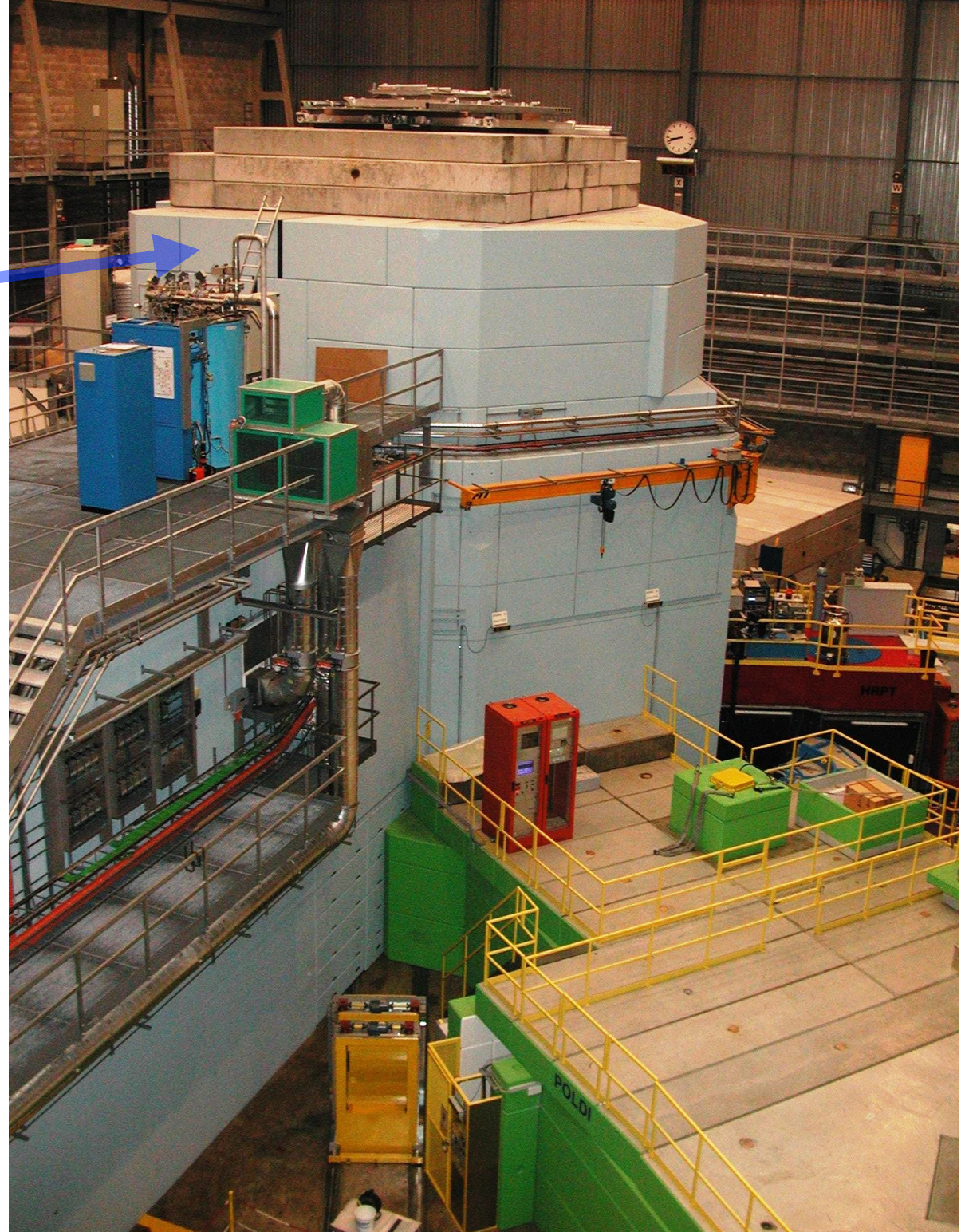


**Some experimental details of neutron  
diffraction experiment at HRPT/SINQ**



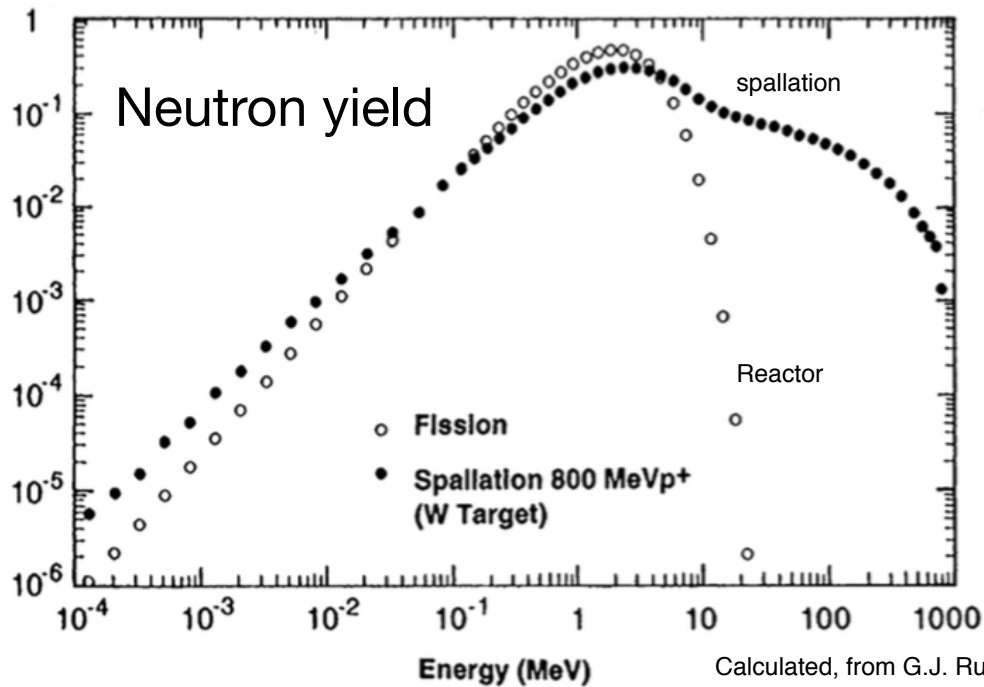
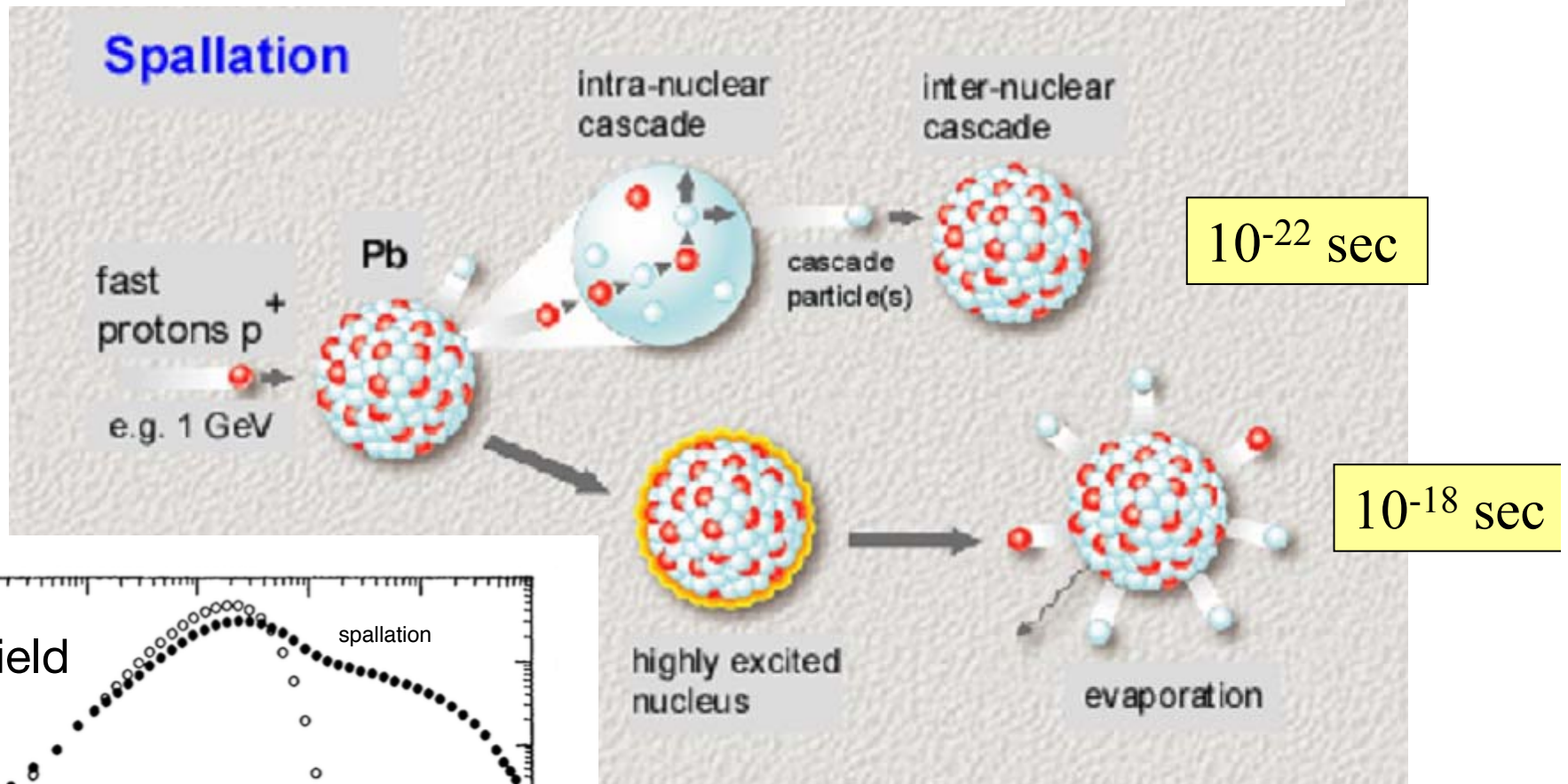
# SINQ hall

The spallation neutron source **SINQ** is a continuous source - the first and the only of its kind in the world - with a **flux of about  $10^{14}$  n/cm<sup>2</sup>/s**. Beside thermal neutrons, a cold moderator of liquid deuterium (cold source) slows neutrons down and shifts their spectrum to lower energies.



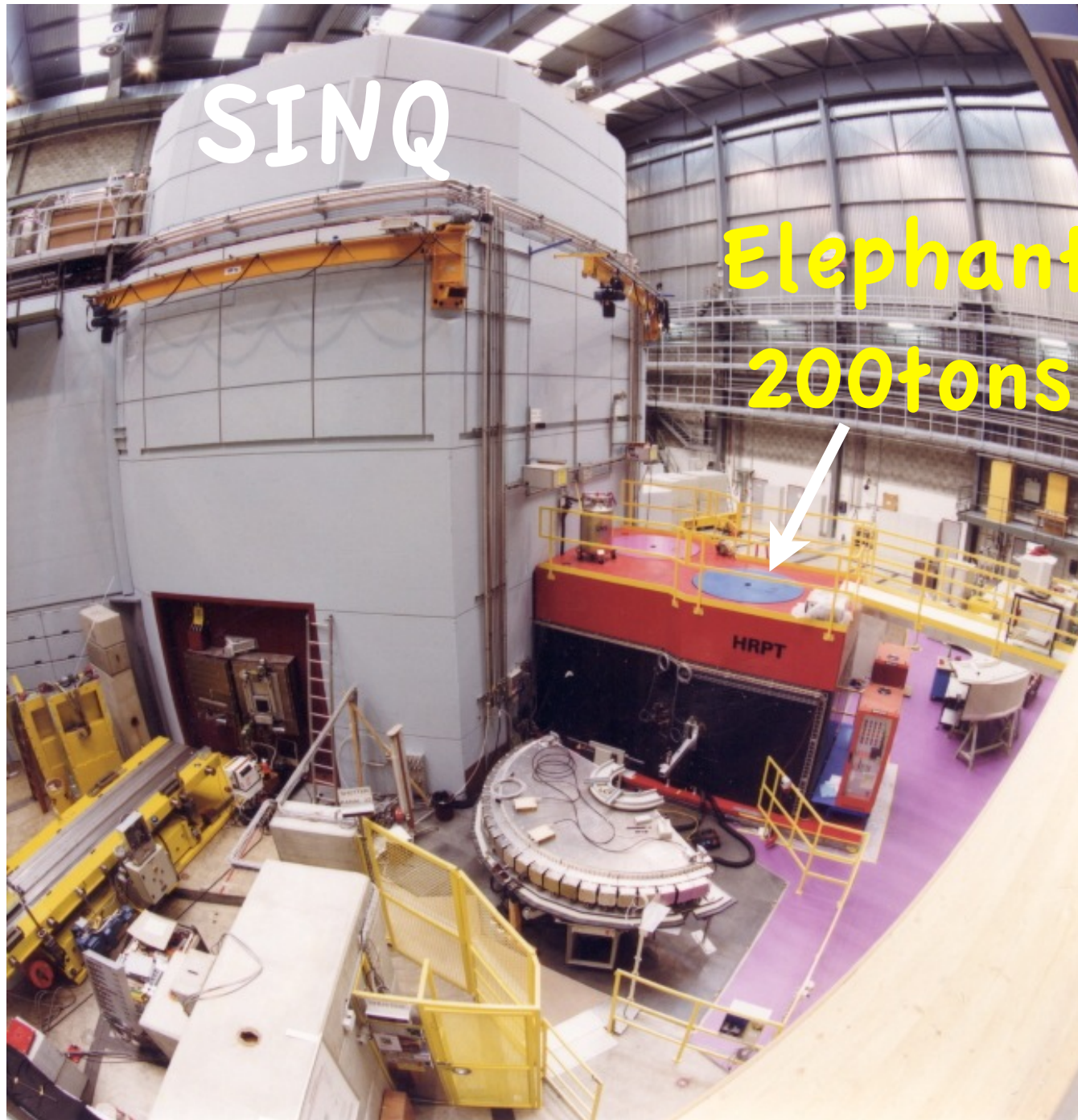


# Spallation, SINQ 590MeV protons, Pb target



Calculated, from G.J. Russell, Spallation physics—an overview, Proceedings of ICANS-XI

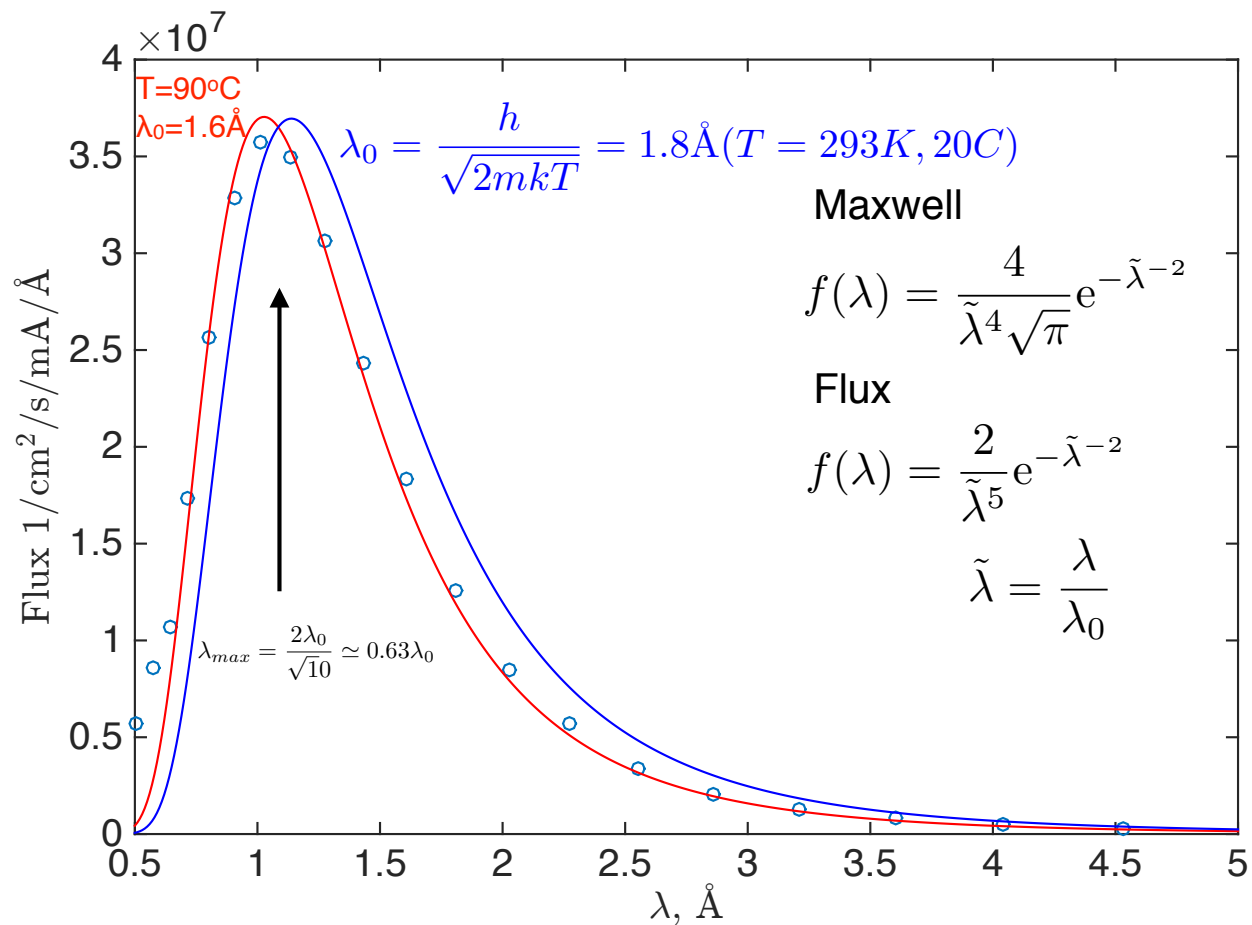
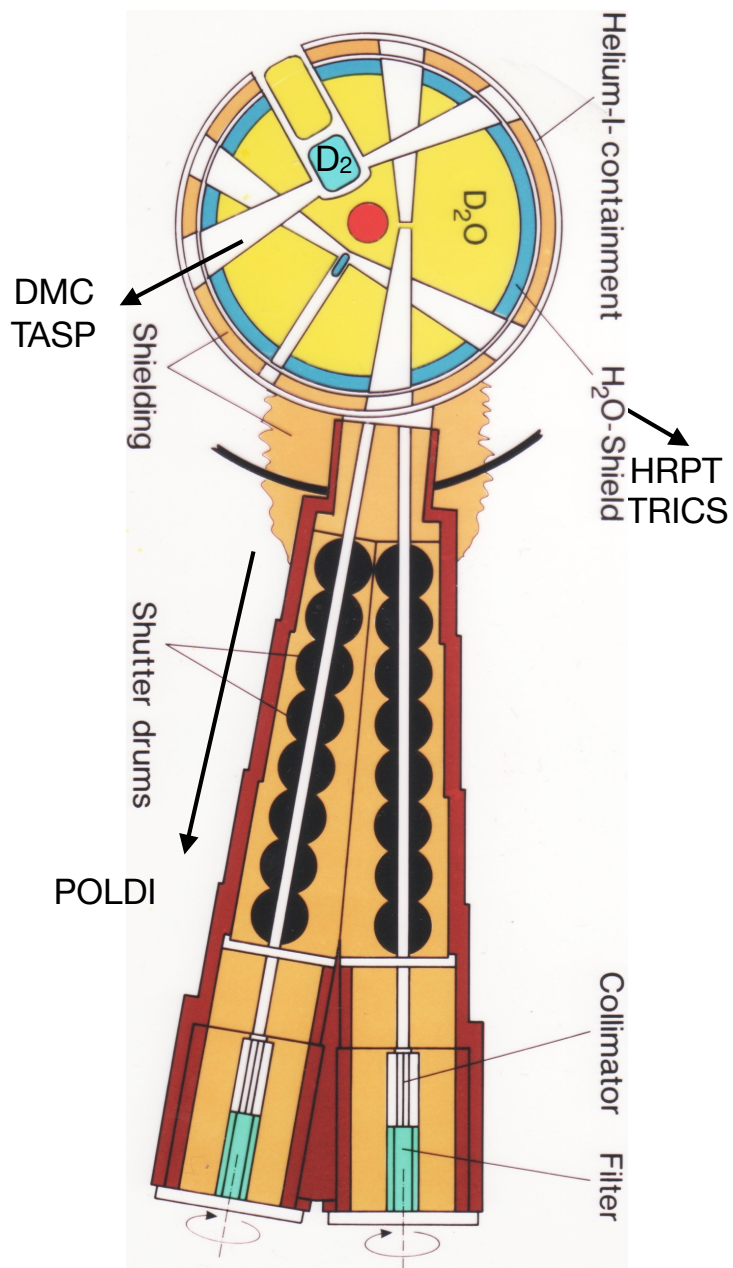
# SINQ hall



Elephant is:  
Shielding of the direct  
neutron beam also from  
fast neutrons for  
diffraction instruments



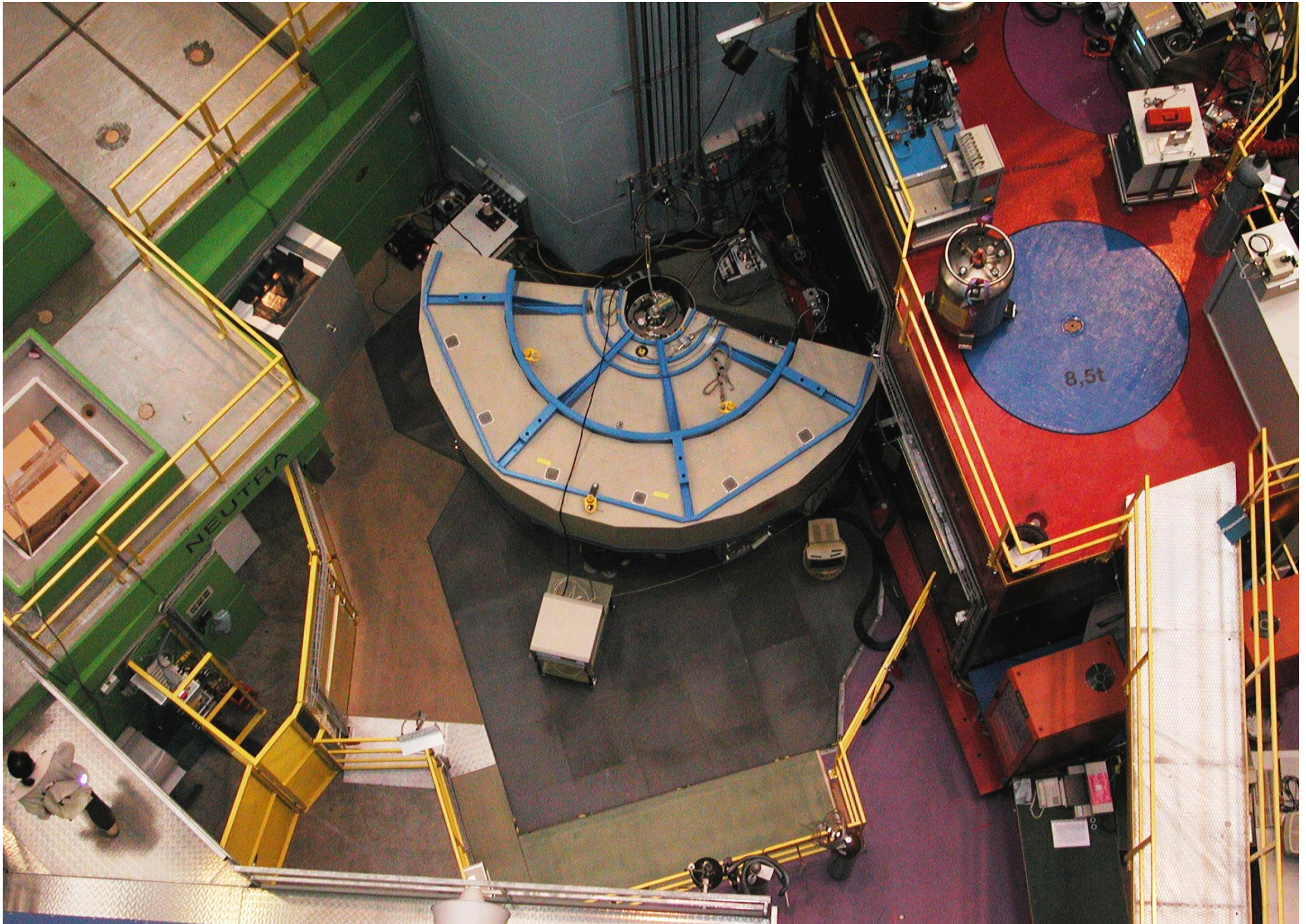
# Neutron (thermal) flux from the D<sub>2</sub>O moderator, Maxwellian at 90°C (HRPT, TRICS)



Total:  $5 \cdot 10^7$   $1/\text{cm}^2/\text{s}/\text{mA}$   
 at SINQ current 2mA:  $10^8$



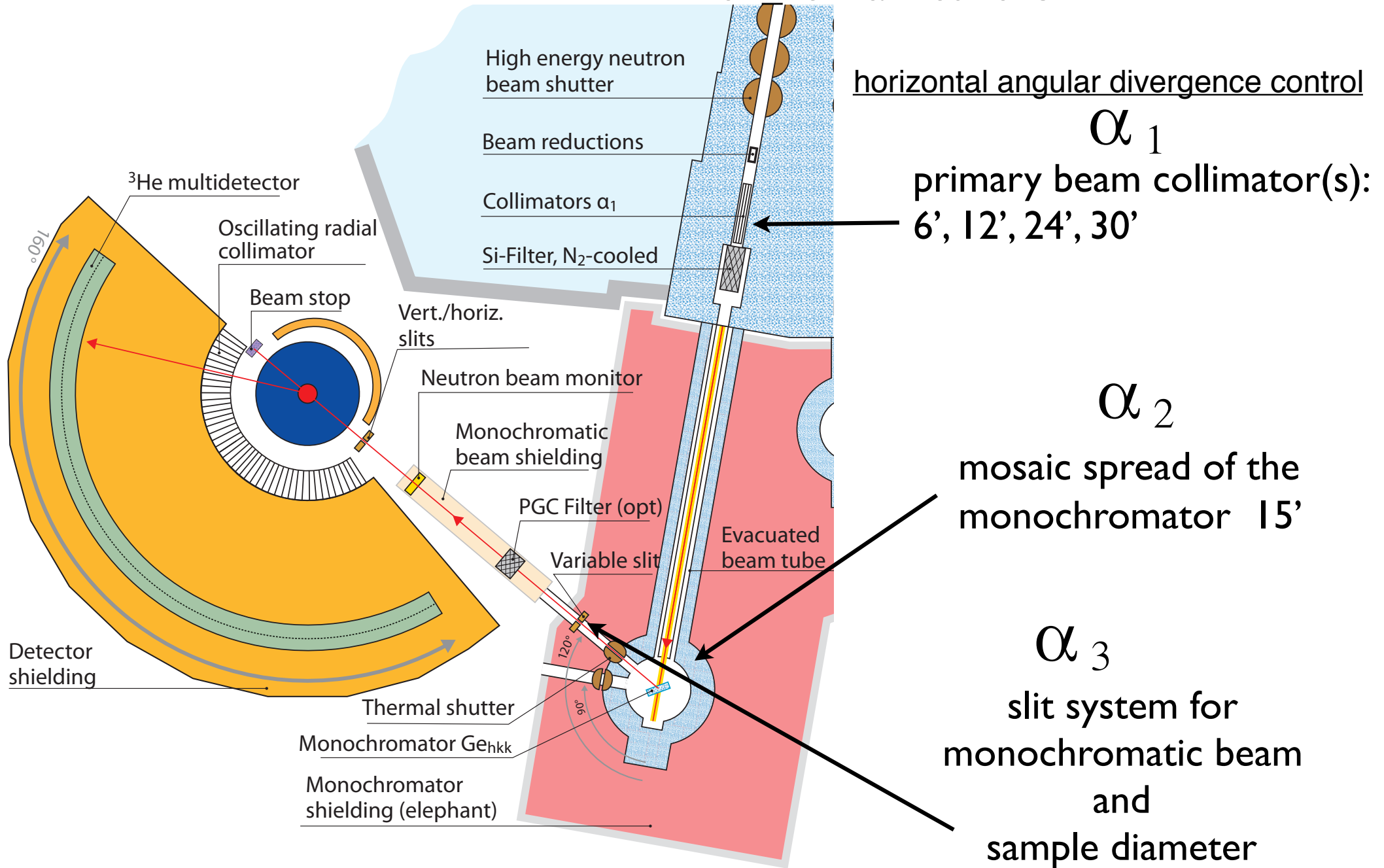
# HRPT areal





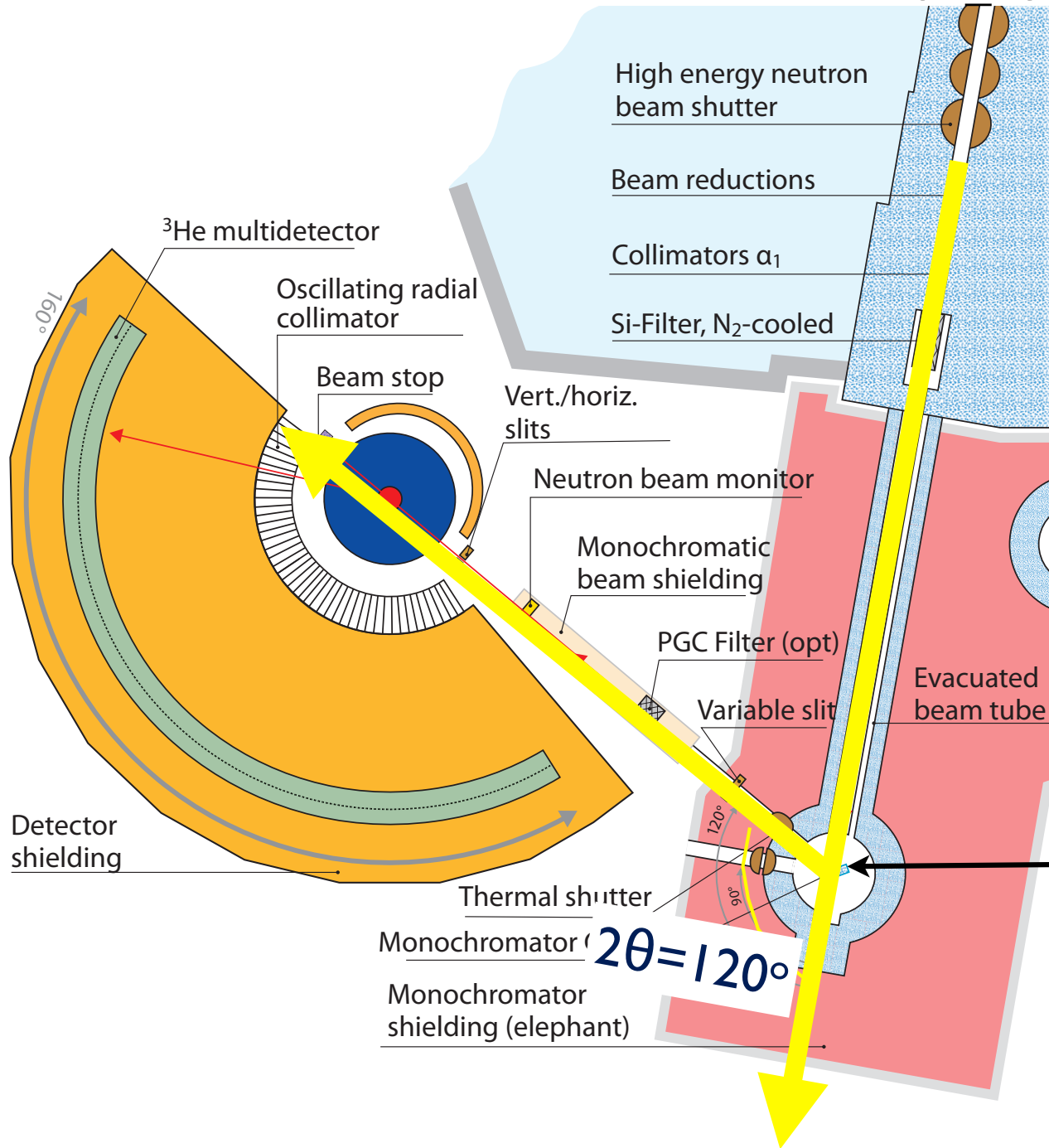
# HRPT layout

## High Resolution Powder Diffractometer for Thermal Neutrons



# HRPT layout

## High Resolution Powder Diffractometer for Thermal Neutrons



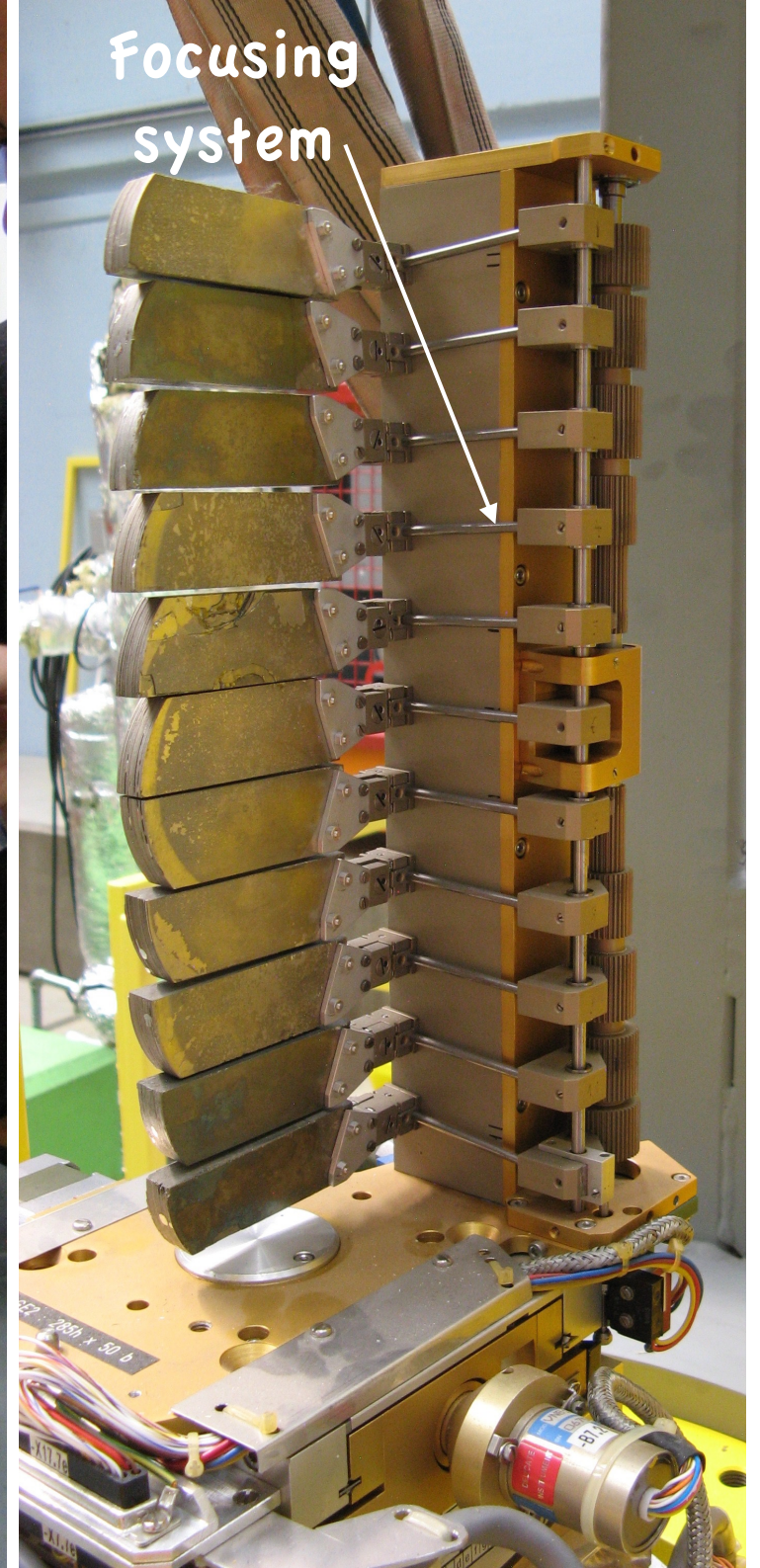
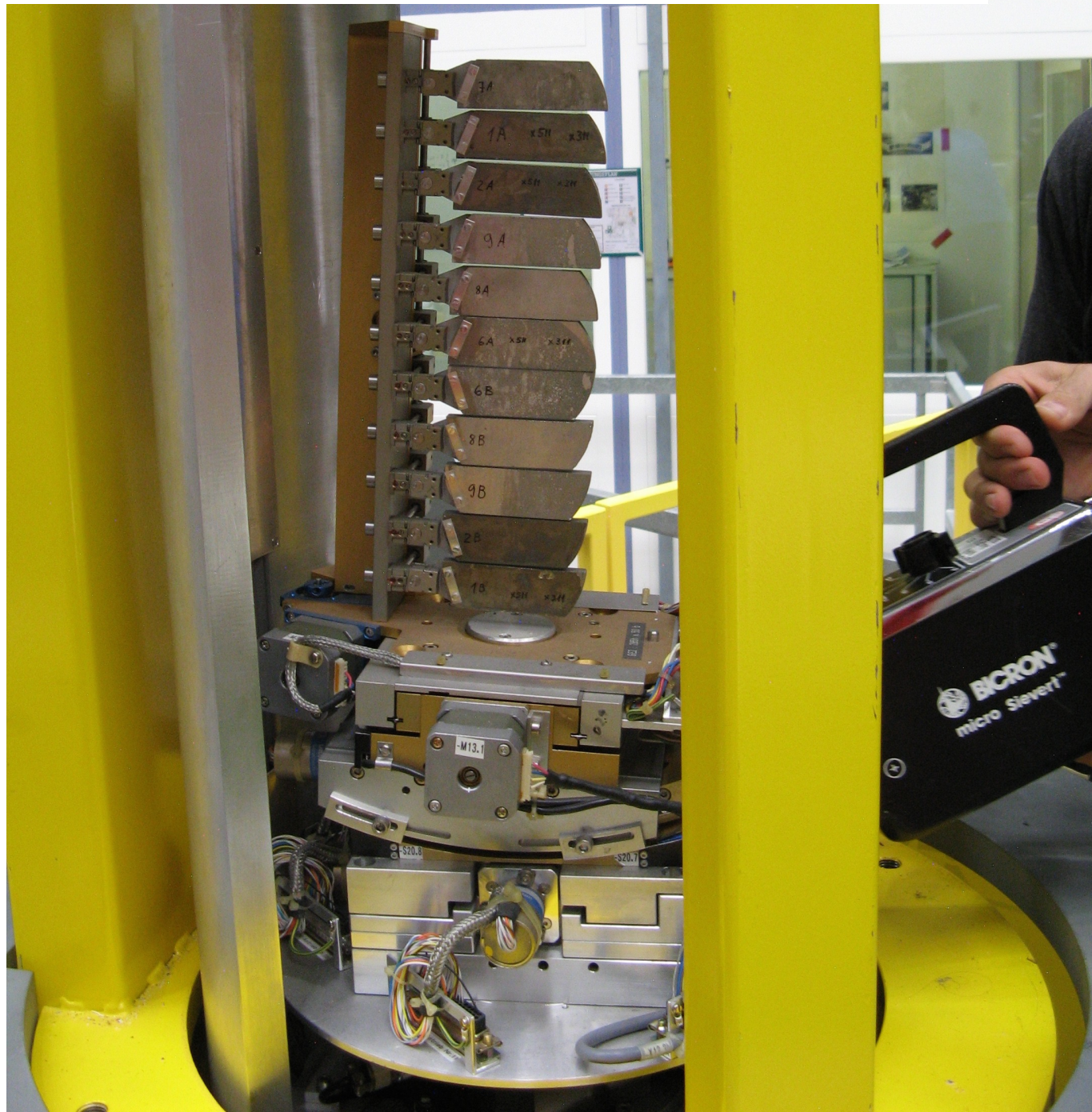
neutron monochromator  
fixed  $120^\circ$  take off angle

$$\lambda = 2d \sin(\theta)$$

$$\lambda = 2d \sin(60^\circ)$$



# Ge single crystal monochromator, 7 motors



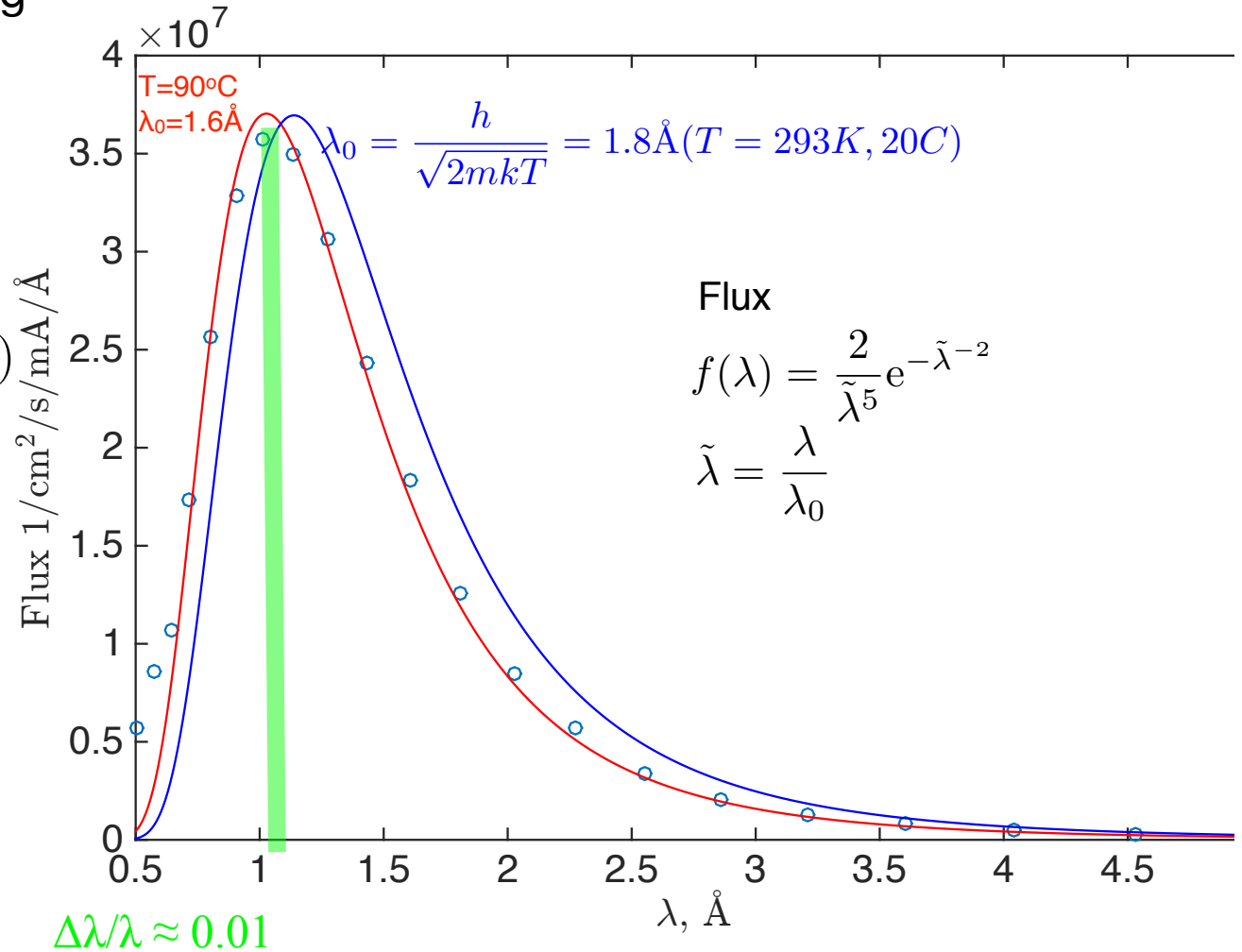


# Monochromator cuts narrow wavelength range from the “white” flux. HRPT $\lambda=0.94 - 2.96 \text{ \AA}$

Intensity of Bragg scattering from big single crystal: Lorentz factor, extinction, geometry, ...

$$I \sim f(\lambda) \Delta\lambda C(\lambda, \theta) \sim f(\lambda) \lambda^{2.5} C'(\theta)$$

for fixed monochromator take-off  $2\theta$  for HRPT



Total:  $5 \cdot 10^7 \text{ 1/cm}^2\text{/s/mA}$   
 at SINQ current 2mA:  $10^8$

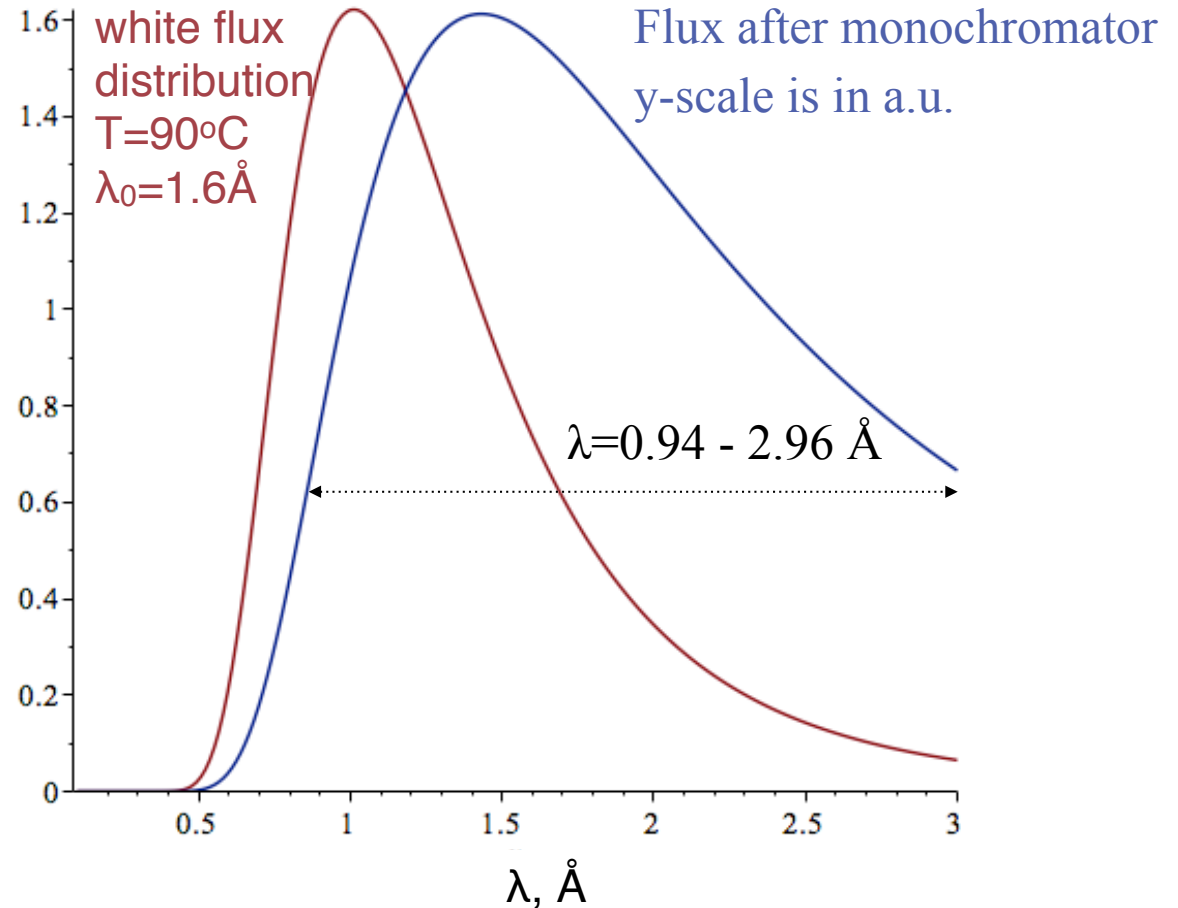
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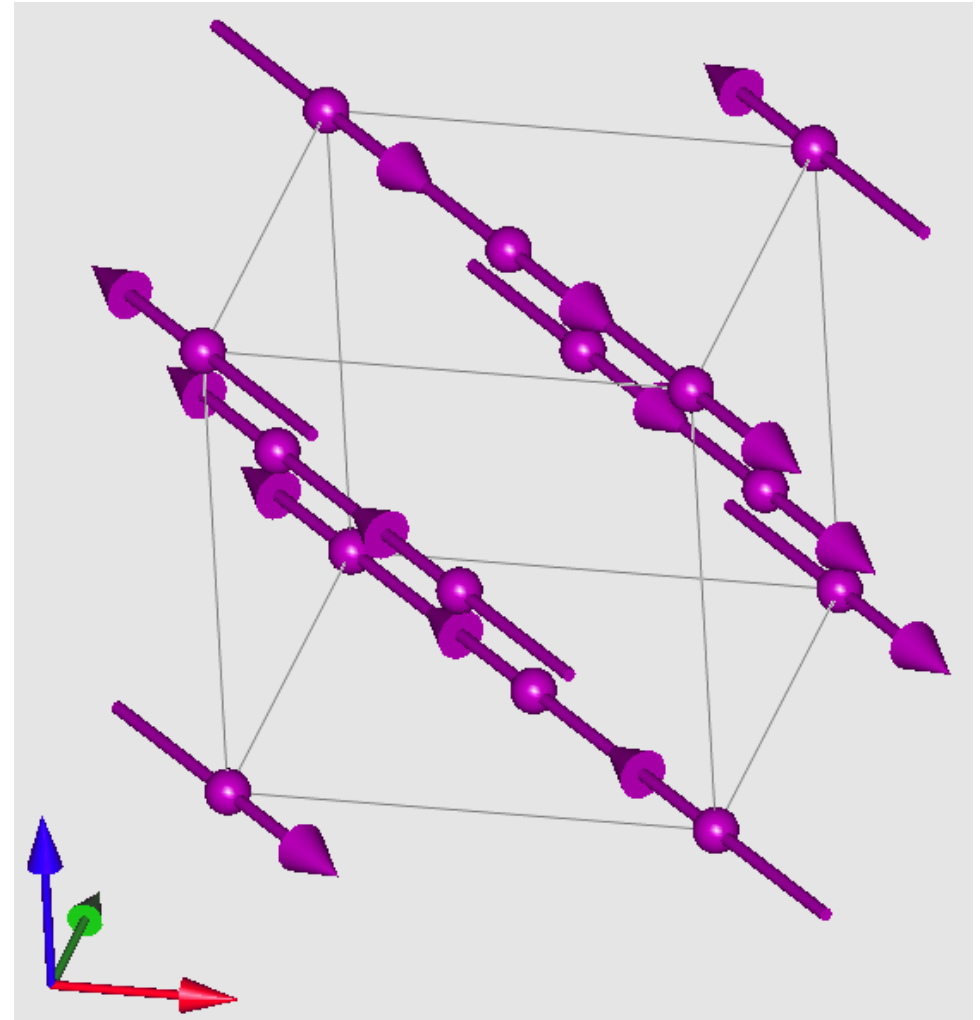
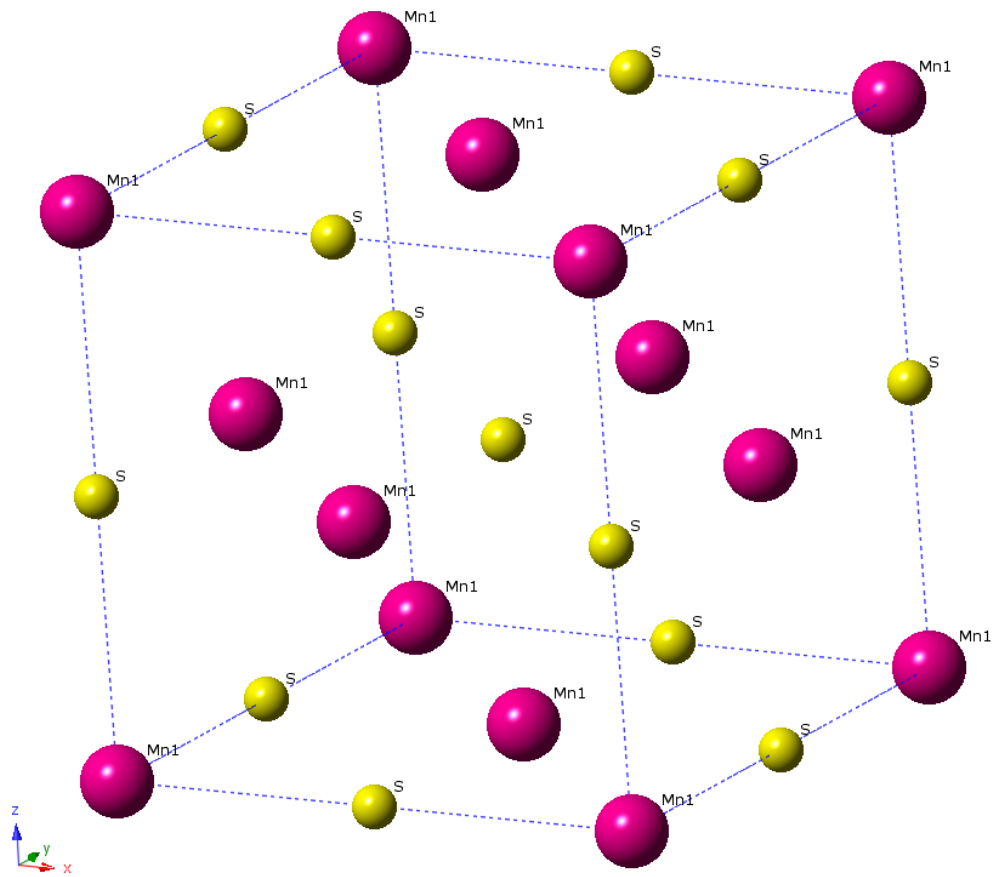
for fixed monochromator take-off  $2\theta$  for HRPT



# Reminder on nuclear and magnetic neutron structure factors



# Approximate crystal and magnetic structures of MnS below Néel temperature



# neutron diffraction experiment ( $\lambda = \text{const}$ )

Nuclear structure factor

$$F(\mathbf{q}) = \sum_j b(\mathbf{r}_j) \exp(i\mathbf{q}\mathbf{r}_j)$$

Magnetic structure factor

$$\mathbf{F}(\mathbf{q}) \propto \sum_j \mathbf{S}_{0\perp j} \cdot \exp(i\mathbf{q}\mathbf{r}_j)$$

atom position and spin

Intensity in the detector

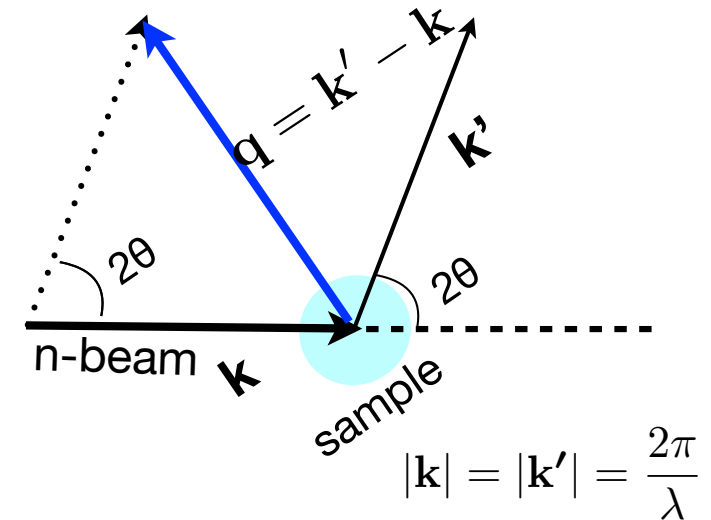
$$\frac{d\sigma}{d\Omega} \propto \mathbf{F}(\mathbf{q})\mathbf{F}^*(\mathbf{q}) \cdot \delta(\mathbf{H} - \mathbf{q})$$

momentum transfer or scattering vector  $\mathbf{q} = \mathbf{k}' - \mathbf{k}$

$$q = \frac{4\pi \sin(\theta)}{\lambda}, \theta = 0.. \leq 90\text{deg.}$$

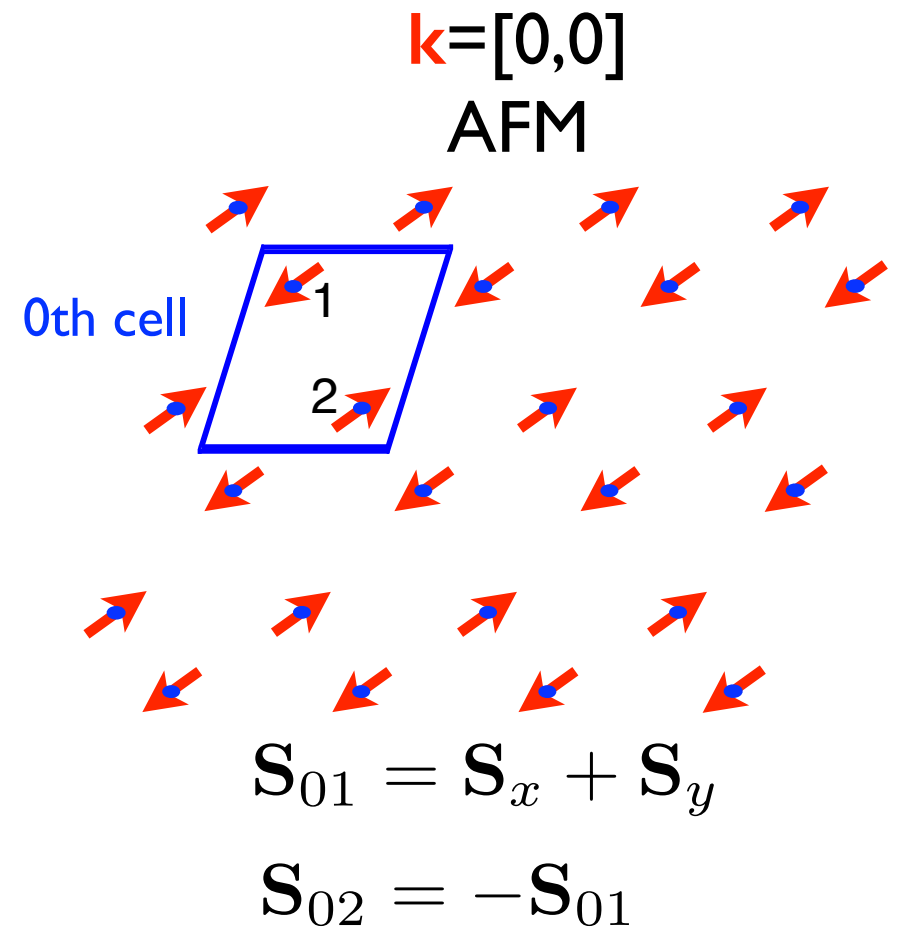
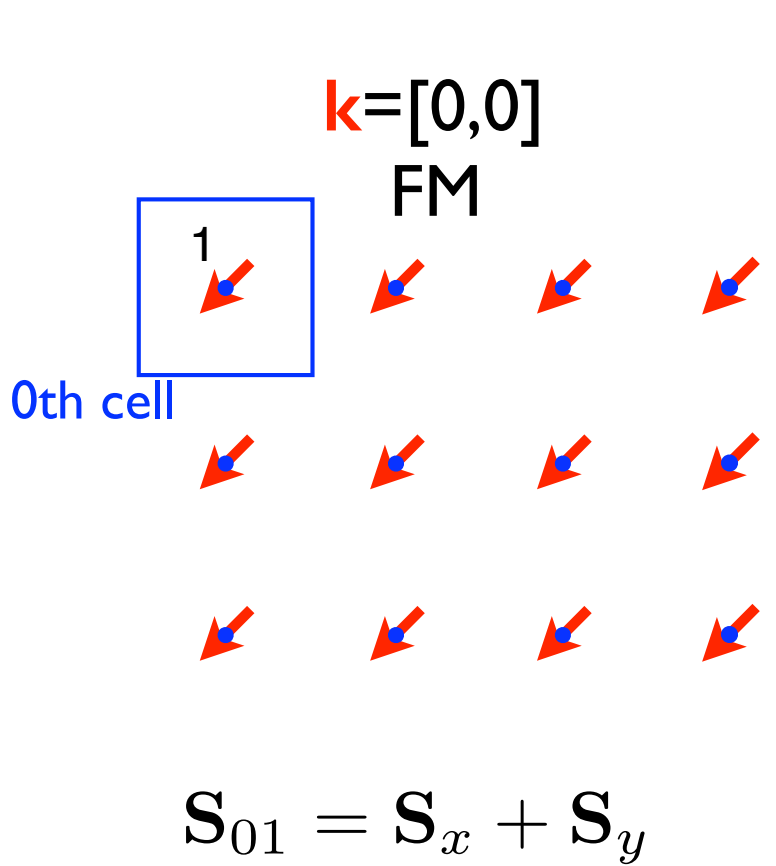
$$q_{max} = \frac{4\pi}{\lambda} \quad \text{large } q_{max} \rightarrow \text{good spatial resolution}$$

( $d_{min} = \frac{\lambda}{2}$ )



# Magnetic structure

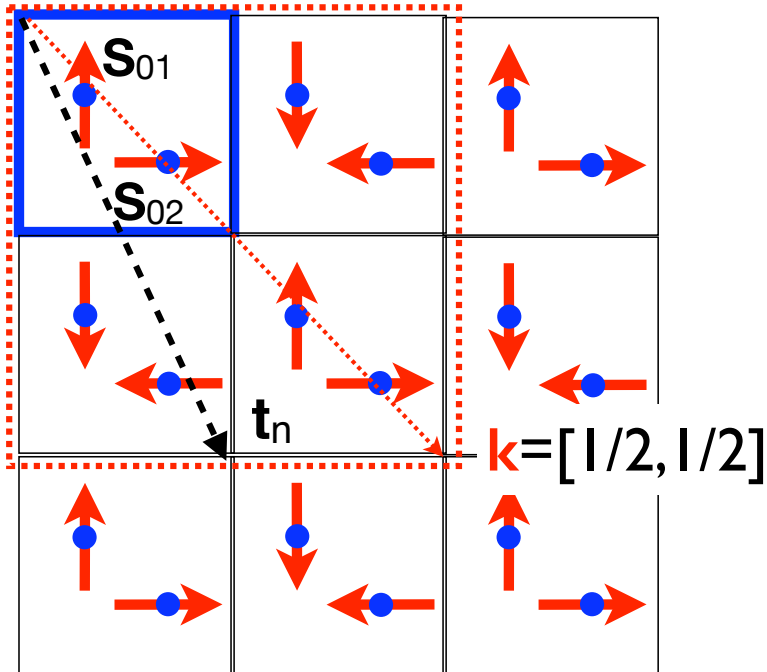
## Examples



Examples of magnetic structures. Propagation vector formalism  $\mathbf{k} \neq 0$ .  
 Magnetic mode  $\mathbf{S}_0$  is specified in zeroth block of the cell == parent cell without centering translations

Magnetic moment  
below a phase transition

0th cell with many atoms in general



$$\mathbf{S}(\mathbf{t}_n) = \text{Re} \left( C \mathbf{S}_0 e^{2\pi i \mathbf{t}_n \mathbf{k}} \right) \sim \cos(2\pi \mathbf{t}_n \mathbf{k} + \varphi)$$

amplitude or mixing coefficients  $\uparrow$   
 magnetic mode  $\nwarrow$

E.g., atom1  $\mathbf{S}_{01} = \mathbf{e}_y$   
 atom2  $\mathbf{S}_{02} = \mathbf{e}_x$

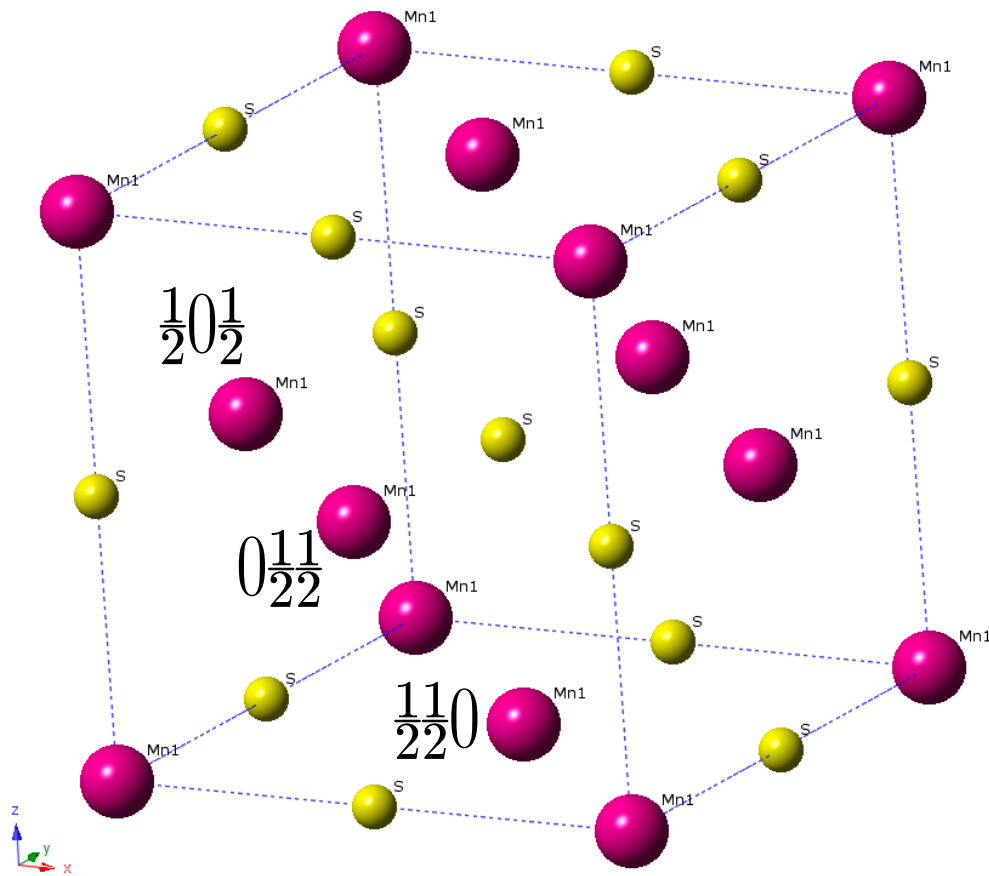
$$\mathbf{S}_1(\mathbf{t}_n) = C \mathbf{e}_y \cos(\pi(t_{nx} + t_{ny}))$$

$$\mathbf{S}_2(\mathbf{t}_n) = C \mathbf{e}_x \cos(\pi(t_{nx} + t_{ny}))$$

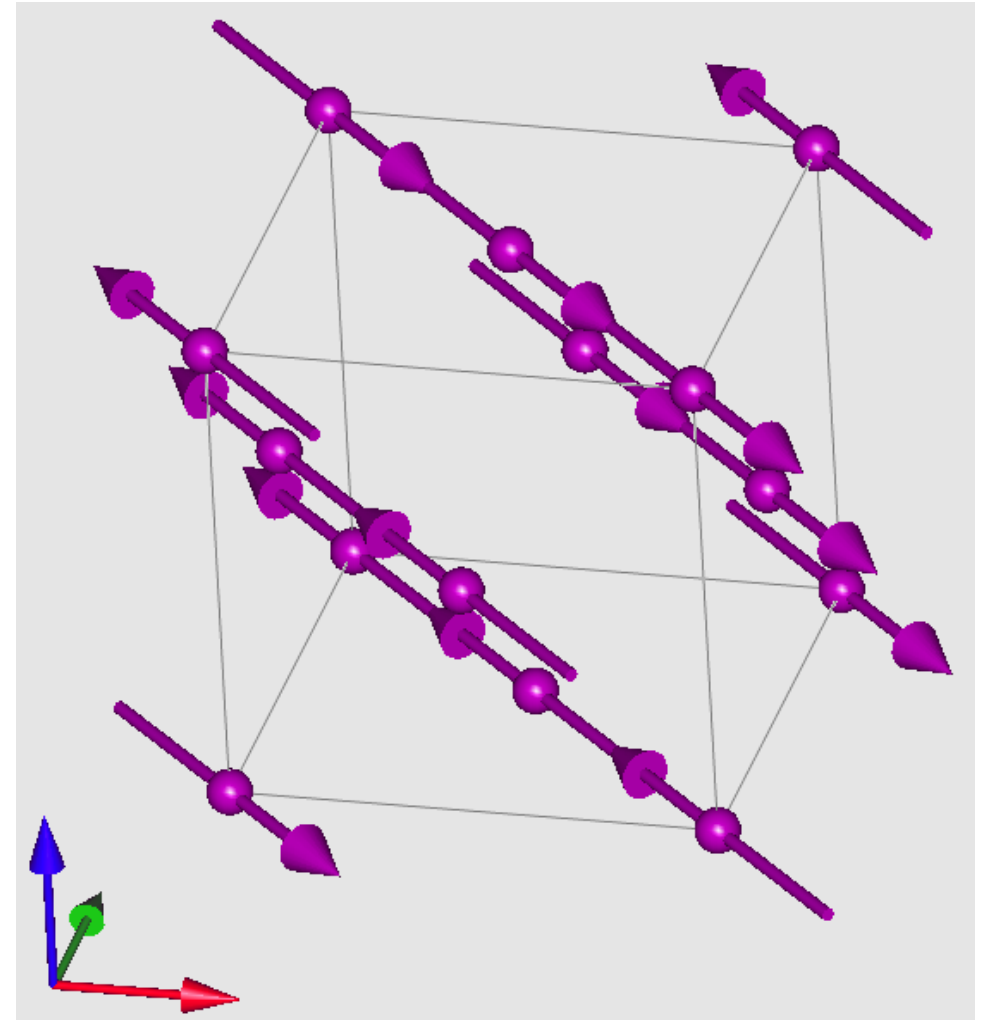
$$t_n = n \cdot 1 \text{ or } n \cdot \frac{1}{2}$$

# Approximate crystal and magnetic structures of MnS below Néel temperature

single propagation vector  
 $\mathbf{k} = [1/2, 1/2, 1/2]$



cubic, Fm-3m: Mn-atom in (000),  
 three other Mn-atoms are generated  
 by F-centering translations



# Scattering from magnetic structure with propagation vector $\mathbf{k}$

In ND experiment we measure correlators of Fourier transform of magnetic lattice

$$\frac{d\sigma}{d\Omega} \propto (\mathbf{F}(\mathbf{q}) \cdot \mathbf{F}^*(\mathbf{q}) + i\mathbf{P} \cdot [\mathbf{F}(\mathbf{q}) \times \mathbf{F}^*(\mathbf{q})]) \cdot \delta(\mathbf{H} \pm \mathbf{k} - \mathbf{q})$$

↑  
structure factor

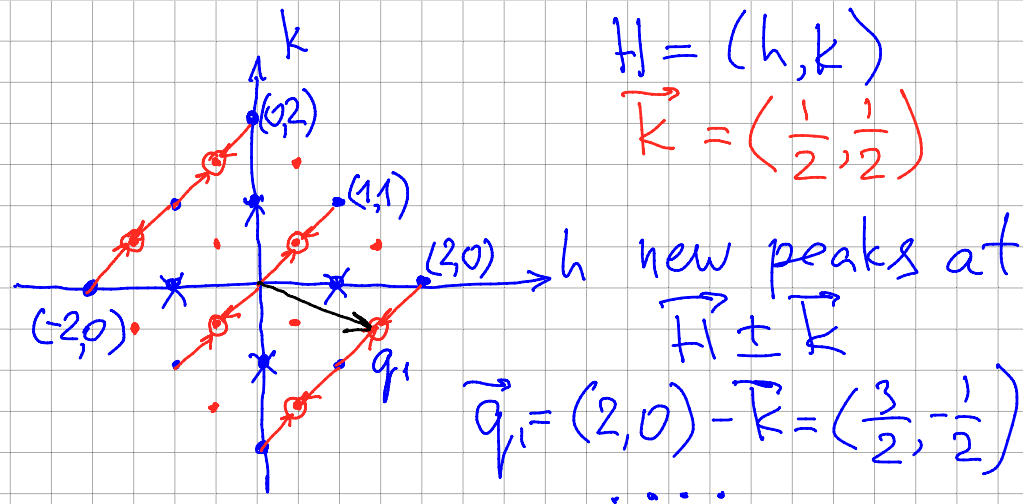
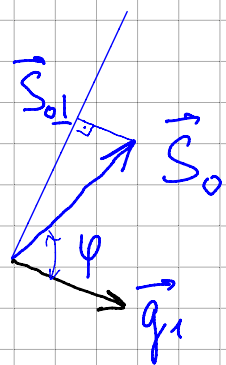
↑  
polarized neutron  
(chiral) term.

↑  
Bragg peak at  
 $\mathbf{q} = \mathbf{H} \mp \mathbf{k}$

$$\mathbf{F}(\mathbf{q}) \propto \sum_j \mathbf{S}_{0\perp j} \cdot \exp(i\mathbf{q}\mathbf{r}_j)$$

$$\vec{S}_\perp = \frac{\vec{q} \times \vec{S} \times \vec{q}}{|\vec{q}|^2}$$

$$|\vec{S}_\perp| = |\vec{S}_0| |\sin\varphi|$$



# Practicum problems

## MAGNETIC ORDER IN MnS

### 5. Practical course at SINQ

#### 5.1 Manganese sulfide MnS

- rock salt crystal structure
- ionic crystal:  $\text{Mn}^{2+}$ ,  $\text{S}^{2-}$
- lattice constant  $a = 5.199 \text{ \AA}$  at  $T = 4.2 \text{ K}$
- space group  $Fm\bar{3}m$
- electronic configuration of  $\text{Mn}^{2+}$ :  $3d^5$
- Néel temperature  $T_N = 161 \text{ K}$
- long-range antiferromagnetic order: antiferromagnetic stacking along (111) of ferromagnetic planes
- therefore doubling of the magnetic unit cell with respect to the crystallographic unit cell



# Task 1: positions of nuclear Bragg peaks, indexing of the peaks

## 5.2 Neutron diffraction of MnS at room temperature

For all measurements at HRPT we will use  $\lambda = 1.886 \text{ \AA}$

$$\lambda = 2d_{hkl} \sin \theta_{hkl} \quad \text{Bragg law} \quad (1)$$

$\lambda$ : neutron wavelength,  $d_{hkl}$ : d-spacing of scattering plane  $hkl$   
 $\theta_{hkl}$ : (half) scattering angle of reflection  $hkl$  in diffraction pattern

$$d_{hkl} = \frac{a}{\sqrt{h^2 + k^2 + l^2}} \quad (2)$$

$a$ : cubic lattice constant,  $h, k, l$  indices of scattering plane

$$\vec{\tau}_{hkl} = \frac{2\pi}{a}(h, k, l) \quad \vec{\tau}_{hkl} \equiv \vec{H} \quad (3)$$

corresponding vector in reciprocal space: a node of reciprocal lattice

①

### Tasks:

$I(2\theta)$

$> T_N \approx 150 \text{ K}$

- measure a diffraction pattern of MnS at  $T = 300 \text{ K}$  in the paramagnetic state
- determine peak positions  $\theta$ ,  $d$ -spacings and indices  $(h, k, l)$  for all observed peaks   
 *first 4*

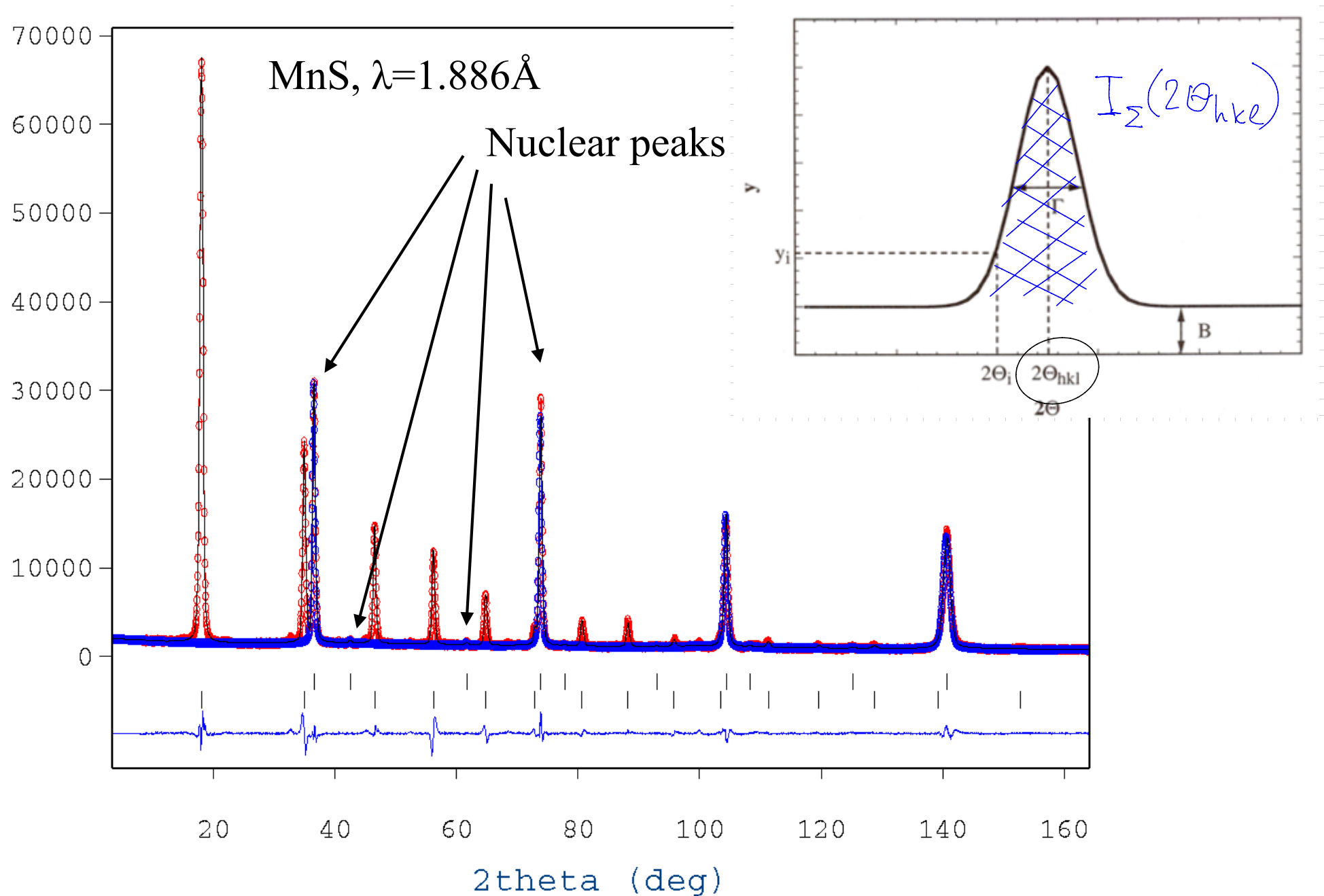
Coherent elastic cross section for nuclear neutron diffraction:

$$\frac{d\sigma}{d\Omega} \sim \sum_{\vec{\tau}_{hkl}} |F_{\vec{\tau}_{hkl}}|^2 \delta(\vec{Q} - \vec{\tau}_{hkl}) \quad (4) \quad \vec{Q} = \vec{\tau}_{hkl} \quad |\vec{Q}| = \frac{4\pi \sin \theta}{\lambda}$$

$\vec{Q}$ : scattering vector,  $\vec{\tau}_{hkl}$ : reciprocal lattice vector defining scattering planes

$\delta(\vec{Q} - \vec{\tau}_{hkl}) \rightarrow$  peak position given by crystal lattice (unit cell)

All the calculations/fits of experimental integrated intensities and peak positions will be done with 'fit' program under HRPT linux-computer



# Task 2a: Calculation of structure factors and Bragg peak intensities and comparison with experiment

Coherent elastic cross section for nuclear neutron diffraction:

$$\frac{d\sigma}{d\Omega} \sim \sum_{\vec{\tau}_{hkl}} |F_{\vec{\tau}_{hkl}}|^2 \delta(\vec{Q} - \vec{\tau}_{hkl}) \quad (4) \quad \vec{Q} = \vec{\tau}_{hkl} \quad |\vec{Q}| = \frac{4\pi\lambda \sin\theta}{\lambda}$$

$\vec{Q}$ : scattering vector,  $\vec{\tau}_{hkl}$ : reciprocal lattice vector defining scattering planes

$\delta(\vec{Q} - \vec{\tau}_{hkl}) \rightarrow$  **peak position given by crystal lattice (unit cell)**

$F_{\vec{\tau}_{hkl}}$ : structure factor of unit cell

$$F_{\vec{\tau}_{hkl}} = \sum_{\vec{d}_i} \left( b_{\vec{d}_i} e^{i\vec{\tau}_{hkl} \cdot \vec{d}_i} \cdot e^{-B_i (Q/4\pi)^2} \right), \text{ we assume } B=0, \text{ but actually } B \approx 0.8 \text{ \AA}^2 \text{ @ RT} \quad (5)$$

$\vec{d}_i$ : atomic coordinate of i-th atom in real space, sum runs over all atoms in unit cell

$b_{\vec{d}_i}$ : scattering length of atom at position  $\vec{d}_i$

Intensity  $\sim |F_{\vec{\tau}_{hkl}}|^2 \rightarrow$  **peak intensity is mainly given by arrangement of atoms in unit cell**

The sum runs over all atoms in unit cell.

For MnS:  $b_{Mn} = -3.73 \text{ fm}$ ,  $b_S = 2.85 \text{ fm}$

femto or fermi  
 $1 \text{ fm} = 10^{-13} \text{ cm} = 10^{-15} \text{ m}$

$\vec{d}$ -vectors:

Mn:  $\vec{d}_1 = a(0, 0, 0)$   $\vec{d}_2 = a(0, \frac{1}{2}, \frac{1}{2})$   
 $\vec{d}_3 = a(\frac{1}{2}, 0, \frac{1}{2})$   $\vec{d}_4 = a(\frac{1}{2}, \frac{1}{2}, 0)$   
 S:  $\vec{d}_5 = a(\frac{1}{2}, \frac{1}{2}, \frac{1}{2})$   $\vec{d}_6 = a(\frac{1}{2}, 0, 0)$   
 $\vec{d}_7 = a(0, \frac{1}{2}, 0)$   $\vec{d}_8 = a(0, 0, \frac{1}{2})$

F-centering. (8)

## Task 2b: Calculation of structure factors and Bragg peak intensities and comparison with experiment

$I_{\Sigma}(2\theta)$  is measured in the experiment  $|\vec{Q}| = \frac{4\pi\lambda\sin\theta}{\lambda}$

For cylindrical geometry of the powder sample container the integrated intensity of the scattered neutrons of the Bragg peak at  $|\vec{Q}|$  is given by

$$I(Q) = C \cdot A(\theta) \cdot L(\theta) \cdot \frac{d\sigma}{d\Omega} = C \cdot A(\theta) \cdot L(\theta) \cdot |F(Q)|^2 \cdot \text{mult}$$

assume:  $A(\theta) \approx 1$  (6)

C: scale factor,  $A(\theta)$ : absorption factor,  $L(\theta)$ : Lorentz factor, *mult*: multiplicity

$$L(\theta) = \frac{1}{\sin\theta \sin 2\theta}$$

(7)

The Lorentz factor  $L(\theta)$  is a geometrical correction depending on the scattering geometry.

2

### Tasks:

- calculate  $|F_{\vec{\tau}_{hkl}}|^2$  for  $(h, k, l) = (1, 1, 1)$ , calculate the multiplicity
- calculate  $|F_{\vec{\tau}_{hkl}}|^2$  for  $(h, k, l) = (2, 0, 0)$ , calculate the multiplicity
- compare ratio of  $I(\vec{\tau}_{111}) / I(\vec{\tau}_{200})$  with the measured intensity ratio

$\rightarrow (111), (-111), \dots$

# Task 3: Indexing of the magnetic Bragg peaks. Calculation of magnetic structure factors and determination of the value and direction of the Mn-spins.

## 5.3 Neutron diffraction of MnS in the magnetically ordered state

3

### Tasks:

- measure a neutron diffraction pattern of MnS at  $T = 50$  K in the magnetically ordered state
- compare this data with the paramagnetic pattern at room temperature
- index the magnetic peaks, i.e. find  $(h, k, l)$  for each magnetic peak
- based on the indices, what is the magnetic unit cell compared to the crystallographic one

Using eq. (6) for  $(111)$ -peak and eq. (10) for  $(\frac{1}{2}\frac{1}{2}\frac{1}{2})$  magnetic peak determine  $\vec{\mu}$ .

Consider  $\vec{\mu} \parallel (111)$  in calculations

- Which direction of  $\vec{\mu}$  can be excluded?

# Task 3: Indexing of the magnetic Bragg peaks. Calculation of magnetic structure factors and determination of the value and direction of the Mn-spins.

Coherent elastic cross section for antiferromagnetic order

$$\frac{d\sigma}{d\Omega} \sim \sum_{\vec{r}_{M,hkl}} |F_{M,hkl}|^2 \delta(\vec{Q} - \vec{r}_{M,hkl})$$

$F_{M,hkl}$ : antiferromagnetic structure factor

Formula (10) is actually formula (6) with  $|F_{M\perp}|$  instead of  $|F|^2$  with the same scale factor  $C$

The intensity of the magnetic Bragg peak at  $|\vec{Q}_M|$  is

$$I(Q_M) = C \cdot A(\theta) \cdot L(\theta) \cdot |F_{M\perp}|^2 \cdot \text{mult}$$

where

$$\vec{F}_{M\perp} = \frac{1}{2} r_0 \sum_{j=1}^4 e^{i\vec{Q}_M \cdot \vec{d}_j} \vec{\mu}_{j\perp} \quad (11) \quad \text{and} \quad \vec{\mu}_{\perp} = \left( \vec{\mu} - \frac{\vec{Q}_M (\vec{\mu} \cdot \vec{Q}_M)}{Q_M^2} \right) \quad (12)$$

where  $\vec{\mu}$  is the magnetic moment in units  $\mu_B$  and  $r_0 = -0.54 \cdot 10^{-12}$  cm,  $\vec{Q}_M \equiv \vec{r}_{M,hkl}$

The sum runs over Mn-atoms.

$\vec{\mu}_1 \uparrow \downarrow \vec{\mu}_{2,3,4}$  Can you explain why?